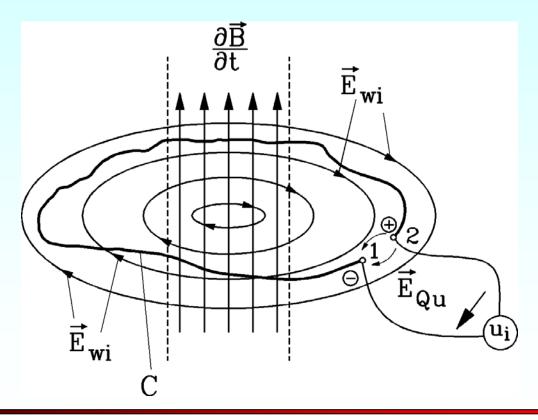
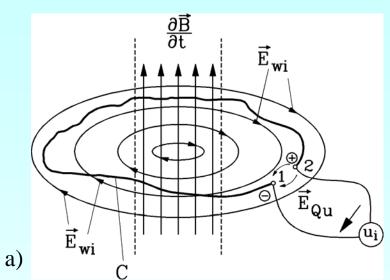
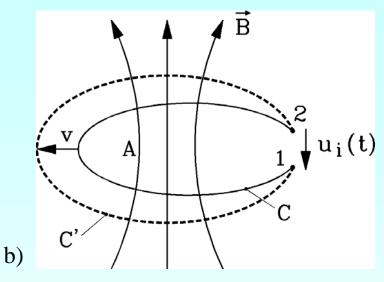
4. Voltage Induction in Three-Phase Machines





FARADAY's law of induction





Each change of flux Φ , which is linked to conductor loop C, causes an induced voltage u_i in that loop; the induced voltage is the negative rate of change of the linked flux.

 $u_i = -d\Phi / dt$ Fluß: $\Phi = \int \vec{B} \cdot d\vec{A}$

- If coil is used instead of loop with N series connected turns, so u_i is N-times bigger: $u_i = -N \cdot d\Phi / dt$
- "Flux linkage" $\Psi = N \cdot \Phi \implies u_i = -d\Psi / dt$
- Changing of Ψ : a) B is changing, b) Area A is changing with velocity v





Induction in resting and moving coils

Resting coils	Moving coils	
Flux density B is changing with time	Flux density B is constant with time	
Coil at rest	Coil moving with velocity v	
$u_i = -d\Psi/dt = -N \cdot d\Phi/dt$		
$u_i = -\partial \Psi / \partial t = \oint \vec{E}_{Wi} \cdot d\vec{s}$	$u_i = \oint (\vec{v} \times \vec{B}) \cdot d\vec{s} = \oint \vec{E}_b \cdot d\vec{s}$	
Electric field strength \vec{E}_{Wi} $(\vec{E}_{Wi} \Leftrightarrow -\partial \vec{B}/\partial t)$	Electric field strength $\vec{E}_b = \vec{v} \times \vec{B}$	
Application of FARADAY's law:		
Transformer coilsStator coils of AC machines	Rotating armature of DC machines	
Transformer induction	Rotating induction	

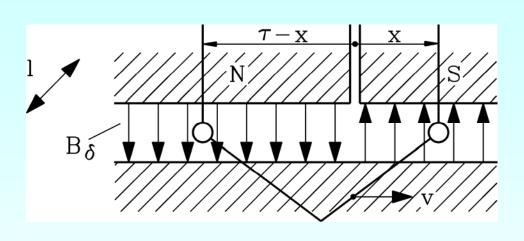
•
$$\frac{d\Phi}{dt} = \frac{d}{dt} \int_{A} \vec{B} \cdot d\vec{A} = \int_{A=const.} \frac{\partial \vec{B}}{\partial t} \cdot d\vec{A} - \oint_{C} (\vec{v} \times \vec{B}) \cdot d\vec{s} \quad \text{(Derivative of product !)}$$

$$u_{i} = \oint_{N \cdot C} (\vec{E}_{Wi} + \vec{E}_{b}) \cdot d\vec{s} = N \cdot \int_{A} -\frac{\partial \vec{B}}{\partial t} \cdot d\vec{A} + N \cdot \oint_{C} (\vec{v} \times \vec{B}) \cdot d\vec{s} = -\frac{d\Psi}{dt}$$



Example: Induced voltage in simple linear machine

• Coil (number of turns N_c , coil span τ) moves within air gap between iron yoke and permanent magnets (Poles N-S-N-S, Pole width $b_p = \tau$) with velocity v.



a) u_i induced in moving coil:

 $\partial B / \partial t = 0$: no change of flux density. Loop C only considered along length 2I, as winding overhang outside of magnetic field.

 \vec{v} , \vec{B} , \vec{s} perpendicular to each other: $u_i = N_c \cdot 2 \int_0^{\infty} (\vec{v} \times \vec{B}) \cdot d\vec{s} = 2 N_c v B l$

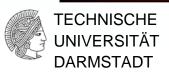
b) $\underline{u_i}$ derived from change of total flux linkage: observer rests with coil: $\underline{u_i} = -d \mathcal{H} dt$: (ALTERNATIVE CALCULATION TO a)!)

Flux linkage changes d Ψ/dt , because coil moves, giving change of coil co-ordinate x = vt!

Coil flux linkage:
$$\Psi = N_c \int \vec{B} \cdot d\vec{A} = N_c \cdot l \cdot [(\tau - x)B_{\delta} - xB_{\delta}] = N_c l B_{\delta} (\tau - 2x)$$

Induced voltage: $u_i = -d\Psi/dt = -N_c l B_\delta \cdot d(\tau - 2 \cdot v \cdot t)/dt = \underline{2N_c v B_\delta l}$

<u>Facit</u>: Induced voltage u_i may be ALWAYS derived from change of total flux linkage.



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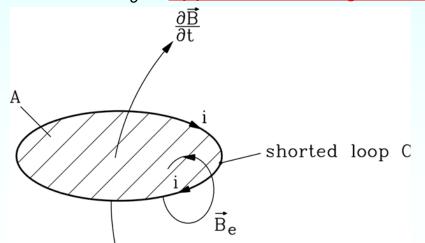
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Law of induction: also called: "LENZ's rule"

Lenz's rule: A change of flux linkage induces voltage u_i , which drives a current i in the loop, which excites a magnetic field B_e , whose direction is opposite to the original change of flux linkage.

- Example: Induction in short circuited loop at rest.
- The change of external field B causes an increase of flux density with orientation from bottom to top. This causes increase of flux in loop area A and induces electrical field E_{Wi} .
- E_{Wi} is left hand oriented to $\partial \vec{B} / \partial t$ and drives in loop C a current i.
- Current i excites (Ampere's law!) a right hand oriented magnetic field B_e.
- Orientation of $B_{\rm e}$ is opposite to change of original flux density $\partial \vec{B} / \partial t$.

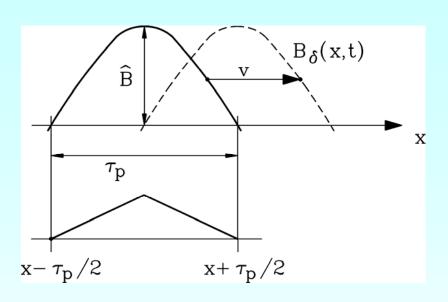


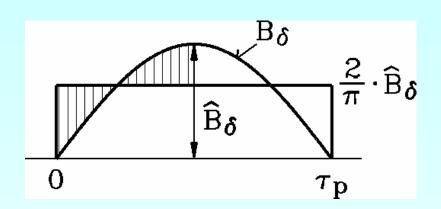
Facit:

The "reaction field" B_e acts AGAINST the original flux density change!



Induction of voltage in stator coil





• Sinusoidal moving wave $B_{\delta 1}(x,t) = \hat{B}_{\delta 1} \cos(x\pi/\tau_p - \omega t)$ causes changing coil flux $\Phi(t)$

$$\Phi(t) = l \int_{-\tau_p/2}^{\tau_p/2} B_{\delta 1}(x, t) dx = \frac{2}{\pi} \tau_p l \hat{B}_{\delta 1} \cdot \cos \omega t \implies \text{flux linkage } \Psi(t) = N_c \Phi(t)$$

• Induced AC voltage in coil is sinusoidal: $u_{i,c}(t) = -d\Psi_c(t)/dt = \hat{U}_{i,c}\sin\omega t$

Voltage amplitude:

$$\hat{U}_{i,c} = \omega N_c \Phi_c = 2\pi f N_c \frac{2}{\pi} \tau_p l \hat{B}_{\delta 1}$$

(full-pitched coil)





Induced voltage by fundamental and harmonic waves

Rotating rotor field (speed n): is a FOURIER-sum of fundamental and harmonic waves:

$$B_{\delta,\mu}(x,t) = \hat{B}_{\delta\mu} \cos(\frac{\mu x \pi}{\tau_{p_{\tau_p}/2}} - \mu \cdot \omega \cdot t), \quad \mu = 1, 3, 5, 7, \dots \quad \omega = 2\pi \cdot n \cdot p$$

• AC coil flux: $\Phi_{c\mu}(t) = l \int_{-\infty}^{\infty} B_{\delta,\mu}(x,t) dx = \frac{2}{\pi} \cdot \frac{\tau_p}{\mu} l \hat{B}_{\delta\mu} \cdot \sin(\frac{\mu\pi}{2}) \cdot \cos(\mu\omega t)$

• Induced voltage:
$$u_{i,c,\mu} = -N_c \frac{d\Phi_{c\mu}}{dt} = \mu\omega \cdot N_c \cdot \frac{2}{\pi} \frac{\tau_p}{\mu} l\hat{B}_{\delta\mu} \cdot \sin(\frac{\mu\pi}{2}) \cdot \sin(\mu\omega t)$$

Facit:

In stator coil not only "useful" voltage due to fundamental (frequency $f = n \cdot p$) is induced, but also harmonic AC voltages with smaller amplitudes, but increased frequencies.

• Smaller voltage amplitudes proportional $\hat{B}_{\delta\mu}$, harmonic frequencies $f_{\mu} = \mu\omega/(2\pi)$.

Note: $\sin(\mu \pi / 2) = (-1)^{(\mu - 1)/2}$ with $\mu = 1, 3, 5, ...$ gives only 1, -1, 1, -1, ... Expression changes only sign, but not amplitude.



Example: No-load voltage in full-pitched coils

• 12-pole synchronous generator: n = 500/min, 2p = 12, full-pitched coils, stator coil data: $N_c = 2$, $W = \tau_p = 0.5$ m, I = 1 m

Fundamental frequency of induced voltage: $f = n \cdot p = (500/60) \cdot 6 = 50$ Hz

ullet Induced harmonic voltage amplitudes depend on rotor air gap field amplitudes $\hat{B}_{_{\!\delta\!\mu}}$

μ	$\hat{B}_{\delta\mu}$	$\hat{B}_{\delta\mu}$ / $\hat{B}_{\delta 1}$	f_{μ}	$arPhi_{C\mu}$	$U_{i,c\mu} = \hat{U}_{i,c\mu} / \sqrt{2}$	$U_{i,c\mu}$ / $U_{i,c1}$
-	Т	%	Hz	mWb	V	%
1	0.9	100	50	286.5	127.2	100
3	0.15	16.7	150	-15.9	-21.2	16.7
5	0.05	5.6	250	3.3	7.1	5.6
7	0.05	5.6	350	-2.3	-7.1	5.6

Facit: Amplitude spectra of inducing field and induced voltage are identical: For a full-pitched coil the spatial field distribution and the time function of voltage are identical!



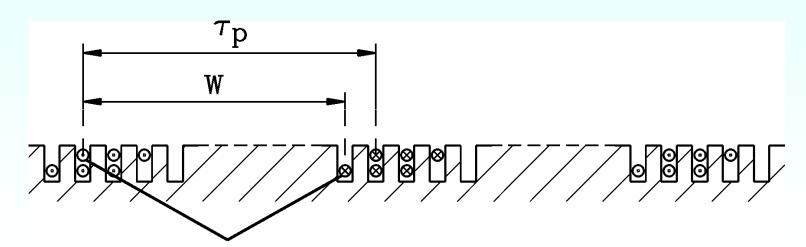
Induction of voltage in pitched coil

• Pitched coil: Coil span is only W instead of τ_p :

$$\Phi_{c\mu}(t) = l \int_{-W/2}^{W/2} \hat{B}_{\delta\mu} \cos(\frac{\mu\pi x}{\tau_p} - \mu\omega t) dx = \frac{2}{\pi} \frac{\tau_p}{\mu} l \hat{B}_{\delta\mu} \cdot \sin(\mu \frac{\pi}{2} \frac{W}{\tau_p}) \cdot \cos\omega t$$

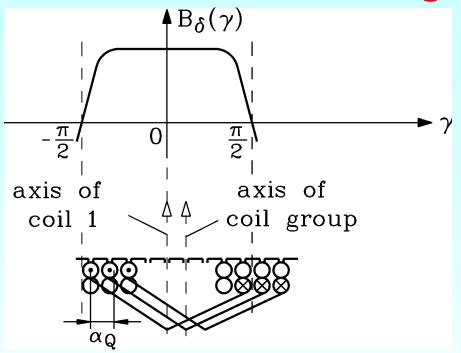
Linked coil flux is smaller by **pitch coefficient** $k_{p,\mu}$, compared to full-pitched coil.

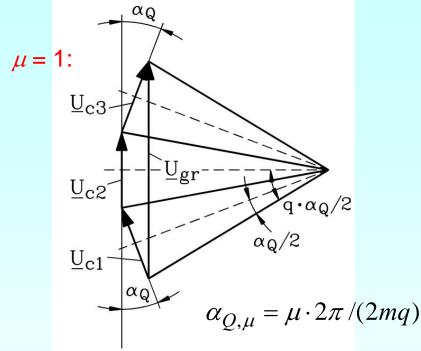
$$k_{p,\mu} = \sin\left(\mu \frac{\pi}{2} \cdot \frac{W}{\tau_p}\right)$$





Induction of voltage in group of coils





•The <u>induced sinusoidal AC voltage per coil group</u> is the sum of complex phasors of the q coils. The coil voltage phasors are phase shifted by angle $\alpha_{Q,\mu}$ between adjacent coils:

• Distribution coefficient:

$$k_{d,\mu} = \frac{\hat{U}_{i,gr,\mu}}{q\hat{U}_{i,c,\mu}} = \frac{2\sin\left(q\frac{\alpha_{Q,\mu}}{2}\right)}{q\cdot 2\sin\left(\frac{\alpha_{Q,\mu}}{2}\right)} = \frac{\sin\left(\mu\frac{\pi}{2m}\right)}{q\cdot \sin\left(\mu\frac{\pi}{2mq}\right)}$$



Induced voltage per phase

- Machine with 2p poles, two-layer winding: One phase consists of 2p coil groups with q pitched coils per group.
- Induced voltage per phase (r.m.s. value):

Fundamental:

$$U_{i1} = \sqrt{2}\pi f \cdot N \cdot k_{w1} \cdot \frac{2}{\pi} \tau_p l \hat{B}_{\delta 1}$$

$$N = 2pqN_c / a \qquad k_{w1} = k_{d1} \cdot k_{p1}$$

$$\mu$$
-th harmonic:

$$U_{i1} = \sqrt{2}\pi f \cdot N \cdot k_{w1} \cdot \frac{2}{\pi} \tau_p l \hat{B}_{\delta 1} \qquad N = 2pqN_c / a \qquad k_{w1} = k_{d1} \cdot k_{p1}$$

$$U_{i,\mu} = \sqrt{2}\pi \mu f \cdot N \cdot k_{w,\mu} \cdot \frac{2}{\pi} \frac{\tau_p}{\mu} l \hat{B}_{\delta \mu}$$

<u>Example:</u> 12-pole synchronous generator: $n = 500/\min$, 2p = 12, f = 50 Hz

- Stator winding: $N_c=2,~q=2,~W=5/6~\tau_p,~a=1,~\tau_p=0.5~\mathrm{m},~l=1~\mathrm{m}$
- Number of turns per phase: $N = 2pqN_c/a = 12 \cdot 2 \cdot 2/1 = 48$

μ	$\hat{B}_{\delta\mu}$	$\hat{B}_{\delta\mu}$ / $\hat{B}_{\delta1}$	f_{μ}	$arPhi_{C\mu}$	$U_{i,\mu}$	$U_{i,\mu} / U_{i,1}$
-	T	%	Hz	mWb	V	%
1	0.9	100	50	276.7	2850.1	100
3	0.15	16.7	150	-11.3	-254.6	8.9
5	0.05	5.6	250	8.0	11.4	0.4
7	0.05	5.6	350	-0.6	-11.4	0.4

Facit: By pitching and by coil group arrangement voltage harmonics are reduced drastically.

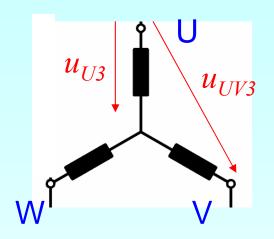


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Star connection: no "third" voltage harmonic



$$u_{U3}(t) = \hat{U}_3 \cdot \cos(3\omega t)$$

$$u_{U3}(t) = \hat{U}_3 \cdot \cos(3\omega t)$$

$$u_{UV3}$$

$$u_{V3}(t) = \hat{U}_3 \cdot \cos(3(\omega t - 2\pi/3)) = \hat{U}_3 \cdot \cos(3\omega t) = u_{U3}(t)$$

$$u_{W3}(t) = \hat{U}_3 \cdot \cos(3(\omega t - 4\pi/3)) = \hat{U}_3 \cdot \cos(3\omega t) = u_{U3}(t)$$
$$u_{UV3}(t) = u_{U3}(t) - u_{V3}(t) = u_{U3}(t) - u_{U3}(t) = 0$$

$$u_{UV3}(t) = u_{U3}(t) - u_{V3}(t) = u_{U3}(t) - u_{U3}(t) = 0$$

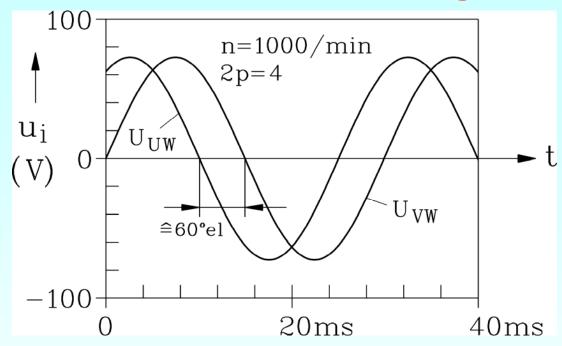
If the stator winding is star connected, the third harmonic voltages in all three phases U, V, W are IN phase and IDENTICAL!

Therefore the <u>line-to-line voltages</u> do not show 3rd harmonic voltage component. Phase voltages in phase cause IN PHASE 3rd harmonic currents, which CANNOT flow at isolated star point (due to 2nd Kirchhoff's law)

$$\underline{I}_3 = \underline{U}_3 / \underline{Z}_3 \implies \underline{I}_{U3} + \underline{I}_{V3} + \underline{I}_{W3} = 3\underline{I}_3 = 0 \implies \underline{I}_3 = 0$$



Star connection: no "third" voltage harmonic



Measured no-load voltage line-to-line of a 4 pole PM synchronous generator at 1000/min, q = 3, skewed slots, star connection, showing nearly ideal sine wave back EMF

Fourier-Analysis of no-load voltage: μ = 1: 33.5 Hz, 74.8 V

 μ = 5: 167 Hz, 0.34 V

Other amplitudes $\mu > 5$ are negligible!



Three phase winding: Self induction leads to magnetizing inductance

Stator air gap field waves, excited by stator current I, induce in stator winding by self induction the voltage u_i!

$$B_{\delta \nu}(x,t) = \hat{B}_{\delta \nu} \cdot \cos\left(\frac{\nu \pi x}{\tau_p} - \omega t\right) \quad \hat{B}_{\delta \nu} = \frac{\mu_0}{\delta} \frac{\sqrt{2}}{\pi} \frac{m}{p} N \frac{k_{w,\nu}}{\nu} I \quad \nu = 1, -5, 7, -11, 13, \dots$$

• Stator air gap field waves $B_{\delta v}(x,t)$: Speed n_v is $n_{\rm syn}/v$. Hence stator field fundamental and field harmonics induce in stator coils **ALL** with the same frequency f.

$$f_{v} = v \cdot p \cdot (n_{syn} / v) = p \cdot n_{syn} = f$$

• r.m.s. of self-induced voltage per phase for each v-th field harmonic:

$$U_{i,\nu} = \sqrt{2}\pi f \cdot N \cdot k_{w,\nu} \cdot \frac{2}{\pi} \frac{\tau_p}{\nu} l\hat{B}_{\delta\nu}$$

• Magnetizing inductance per phase: $L_{h\nu}$ for ν -th air gap field harmonic wave.

$$U_{i,v} = \omega L_{hv} I \quad \Rightarrow \quad L_{hv} = \mu_0 N^2 \frac{k_{w,v}^2}{v^2} \frac{2m}{\pi^2} \frac{l\tau_p}{p \cdot \delta}$$



Stray inductance of stator winding per phase

$$L_{\sigma} = L_{\sigma,Q} + L_{\sigma,b} + L_{\sigma,o}$$

• Air gap field: Fundamental wave = Magnetizing field (subscript h): L_h Magnetizing inductance L_h

$$L_h = L_{h,\nu=1}$$

- Magnetic field in **slots** (slot stray field) and around the **winding overhang** is NOT linked with rotor winding. It does NOT produce any forces with rotor current. Hence it does NOT contribute to electromechanical energy conversion, and is thus called **stray field (subscript σ)**.
- Stray flux induces in stator winding additional voltage by self induction. Hence we define: Slot stray inductance $L_{\sigma Q}$, overhang stray inductance $L_{\sigma b}$: $U_{i\sigma,Q+b} = \omega(L_{\sigma Q} + L_{\sigma b})I$
- Air gap field <u>harmonic</u> waves induce stator winding with voltage $U_{i,v}$ with the same frequency f. So they are summarized as total harmonic voltage : $\sum_{i=1}^{\infty} U_{i,v}$

$$L_{h,total} = \frac{\sum_{\nu=1,-5,7,...}^{\infty} U_{i,\nu}}{\omega I} = \sum_{\nu=1,-5,7,...}^{\infty} L_{h\nu} = (1+\sigma_o)L_{h,\nu=1} \implies \sigma_o = \sum_{\nu=1,-5,7,...}^{\infty} \left(\frac{k_{w,\nu}}{v \cdot k_{w,1}}\right)^2 - 1$$

 σ_o : harmonic stray coefficient (is small: ca. 0.03 ... $\overline{0.09}$).

• Harmonic field waves are linked to rotor, but "disturbe" basic machine function; hence they are summarized in harmonic stray inductance $L_{\sigma o}$: $U_{\underline{i}\sigma,o} = \omega L_{\sigma,o} I$, $L_{\sigma,o} = \sigma_o L_h$

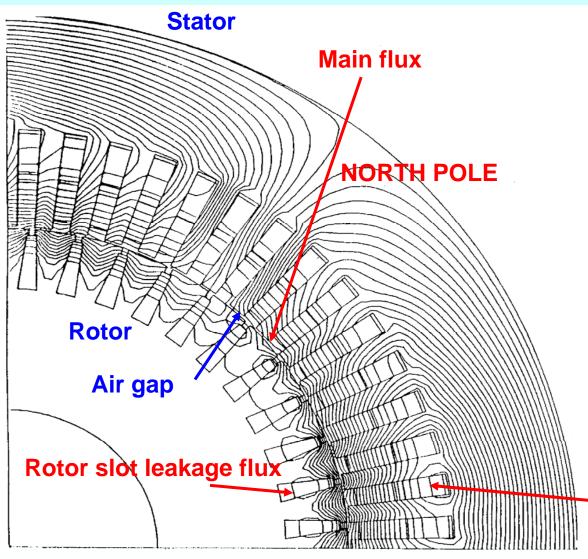


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Field lines B of a cage induction machine



Main flux: Links stator and rotor winding; field lines cross the air gap

Leakage flux (stray flux): Is only linked with either stator or rotor winding; field lines DO NOT cross the air gap

Example:

Four-pole wedge bar rotor: Field lines at stand still (n = 0)

- Rotor frequency = Stator frequency
- Rotor current is NEARLY in phase opposition to stator current

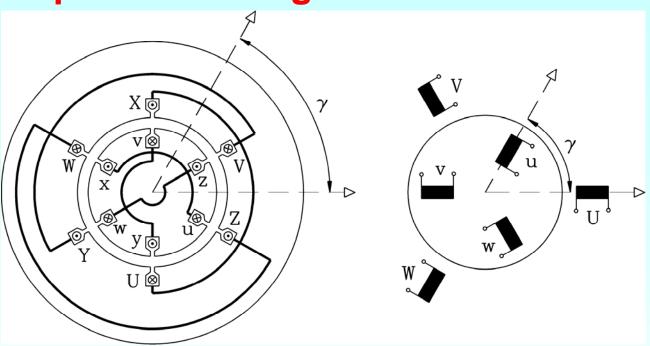
Stator slot leakage flux



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Three phase winding in stator and rotor



- In stator and in rotor each a three-phase winding is arranged:
- in stator: 3 phases between terminals U-X, V-Y, W-Z, subscript s,
- in rotor: 3 phases between terminals u-x, v-y, w-z, subscript r.
- We assume: Rotor is at rest (stand still), and is turned by angle γ with respect to stator. γ = angle between winding axis of stator and rotor winding (= centre of coils).

NOTE: $\gamma = 2\pi$, if rotor is shifted to stator by 2 poles: $2\tau_p$.

Pole numbers of stator and rotor winding must be identical 2p!



Mutual inductance between stator and rotor phase

	Stator	Rotor
Pole count	2р	2p
Phase count	<i>m</i> _s	m_r
Turns/Phase	N_s	N_r
Pitching	$W_{\rm s}/ au_{ m p}$	W_r/ au_p
Coils/group	$q_{\rm s}$	q_r
Slot count	Q_s	Q_r

From now on only fundamental field waves considered!

$$Q_r \neq Q_s$$

• Mutual inductance: e. g.: Stator air gap wave $B_{\delta}(x,t)$ induces voltage in rotor winding:

$$B_{\delta}(x,t) = \hat{B}_{\delta} \cdot \cos(\frac{\pi x}{\tau_p} - \omega_s t) \text{ with amplitudes } \hat{B}_{\delta} = \frac{\mu_0}{\delta} \frac{\sqrt{2}}{\pi} \frac{m_s}{p} N_s k_{ws} I_s$$

• Amplitudes of induced voltages in rotor winding:

$$U_{i,r} = \sqrt{2}\pi f_s \cdot N_r \cdot k_{wr} \cdot \frac{2}{\pi} \tau_p l \hat{B}_{\delta}$$
 Rotor frequency f_r (at locked rotor = stand still): $f_r = f_s$.

• Fundamental wave: Mutual inductance per phase M_{sr} : $U_{i,r} = \omega_s M_{sr} I_s$

$$M_{sr} = \mu_0 N_s k_{w,s} N_r k_{w,r} \frac{2m_s}{\pi^2} \frac{1}{p} \frac{\tau_p l}{\delta}$$

$$M_{sr} = M_{rs}$$

Note:
$$M_{sr} = M_{rs}$$
 at $m_s = m_r$!



Rotary transformer

- Induced rotor voltages are **phase shifted** by angle γ with respect to stator voltages, as rotor is shifted by that angle γ mechanically.
- <u>Series connection of stator and rotor winding</u> U and u (in the same way: V and v; W and w) ⇒ The following resulting voltage occurs between FIRST terminal of stator winding and SECOND terminal of rotor winding (per phase):

$$\underline{U} = \underline{U}_S + \underline{U}_r = U_S + U_r e^{-j\gamma}$$
, e.g. $\underline{U}_r = \underline{U}_s$: $\underline{U} = U_S + U_S e^{-j\gamma} = U_S \cdot (1 + e^{-j\gamma})$

• By turning the rotor we get a continuous change of angle γ .

<u>Facit:</u> With rotary transformer a **continuous change** of output voltage between 0 and $2U_s$ is possible at constant line frequency, which is used in test facilities as variable voltage source.

